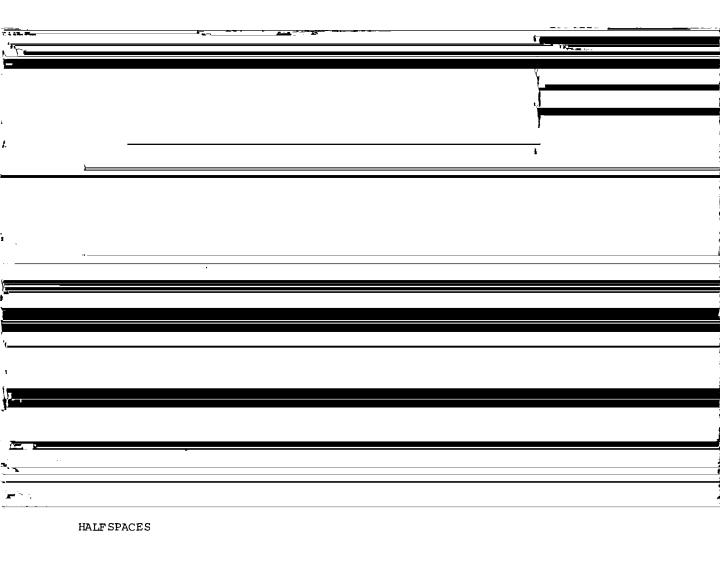
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Introduction

The present work is concerned with the axisymmetric interaction of a Winkler medium contained between two homogeneous, isotropic elastic halfspaces. The Winkler medium, which is composed of a dense array of independent spring

due to the external traction g(r) is given by

where a is a typical length parameter of the problem and G and ν are respectively the linear elastic shear modulus and Poisson's ratio of the

$$\bar{w}_{q}^{0}(\xi) = H\{w_{q}(r);\xi\} = \int_{0}^{\infty} rw_{q}(r)J_{0}(\xi r/a)dr$$
 (2)

and $q^{0}(\xi)$ denotes the zeroth-order Hankel transform of the normal surface traction q(r). The corresponding Hankel inversion theorem is

$$w_{q}(r) = \frac{1}{a^{2}} \int_{0}^{\infty} \xi \overline{w}_{q}^{0}(\xi) J_{0}(\xi r/a) d\xi.$$
 (3)

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shown that the transformed value for the surface displacement $\mathbf{w}_{p}^{\ }(\mathbf{r})$ is given by

$$\frac{-0}{w_p^0}(\xi) = -\frac{a(1-v)}{G\xi} \bar{S}^0(\xi)$$
(4)

where

$$=0$$
, P , ($\xi c/a$), $-\xi c/a$

Compression of The Winkler Layer

	The results described above can now be employed to obtain expressions for the
	transformed values of the surface displacements of the halfspaces under the
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$$\frac{w_{m}(r)}{[ka^{2} (1-v_{m})/G_{m}]} = \int_{0}^{\infty} \frac{\left[\frac{(1-v_{n})}{\xi a^{2}G_{n}} \left\{ \overline{S}_{n}^{0}(\xi) - \overline{S}_{m}^{0}(\xi) \right\} - \overline{S}_{m}^{0}(\xi) \right]}{\left[1 + \frac{ka}{\xi} \left\{ \frac{(1-v_{m})}{G_{m}} + \frac{(1-v_{n})}{G_{n}} \right\} \right]} J_{0}(\xi r/a) d\xi.$$

The specific expressions for the surface displacements of the halfspace regions 1 and 2 are recovered by substituting m=1, n=2 and m=2, n=1 in (10), respectively. In addition, the contact stress q(r) can be directly employed to derive appropriate expressions for the stresses and displacements in the two halfspace regions.

Evaluation of The Infinite Integrals

The general numerical evaluation of the integral expressions for the surface

direct numerical integration technique. Briefly, such numerical integration is performed by representing the integrand as an infinite series bounded by subsequent zeros of $J_0(\xi r/a)$. The application of a Gauss-Legendre quadrature technique for the evaluation of each interval of the integrand yields rapidly

$$+ \frac{(1-\nu_{m})^{k} P_{n}^{c} c_{n}}{4\pi G_{m}^{G} G_{n}^{a}} \left[\left\{ \frac{a}{c_{n}} + (\lambda - \mu_{n}) e^{\lambda c_{n}} \right\}^{a} \operatorname{Ei} \left(-\lambda c_{n}/a \right) \right\} +$$

$$- \left\{ \frac{a}{c_{m}} + (\lambda - \mu_{m}) e^{\lambda c_{m}} \right\}^{a} \operatorname{Ei} \left(-\lambda c_{m}/a \right) \left\{ \frac{1-\nu_{n}}{1-\nu_{m}} \frac{P_{m}^{c} c_{m}}{P_{n}^{c} c_{n}} \right\} \right] .$$
(12)

In (11) and (12) $\left|\arg\lambda\right| < \pi$; $(c_i/a) > 0$ and

$$\mu_{\alpha} = \frac{2(1-\nu_{\alpha})a}{c_{\alpha}} ; \quad \lambda = ka \left[\frac{(1-\nu_{m})}{c_{m}} + \frac{(1-\nu_{n})}{c_{n}} \right] ; \quad (13)$$

also Ei(-x) is the exponential integral, which is related to the function $E_1(x)$ according to Ei(-x) = - $E_1(x)$; tabulated numerical values for $E_1(x)$

case of identical loading by identical halfspaces the equations (11) and (12) reduce to the convenient forms

$$\frac{q(0)}{\sqrt{1+(1-\alpha)}} = \lambda \{\underline{u} - \lambda\} e^{\lambda} Ei(-\lambda) + \{\underline{u} - \lambda + 1\}.$$

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